

# A SIMPLE TUTORIAL ON DISCONTINUOUS GALERKIN METHODS

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*CERMICS, ENPC*

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## Outline

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- Definitions
- A simple example
- Issues
- Historical development
  - elliptic equations
  - hyperbolic equations
- Discontinuous vs. continuous Galerkin
- Further discussion: convergence results, a posteriori results, current research

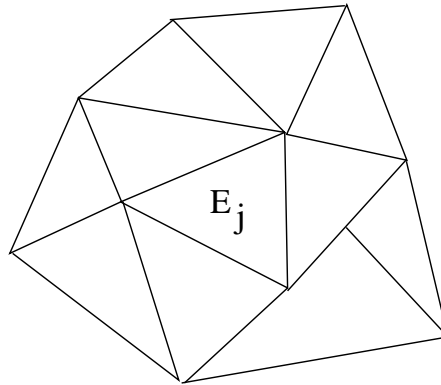
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## Definitions

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- Discontinuous Galerkin methods are those finite element methods which utilize element-wise discontinuous basis functions in the choice of approximation spaces.
- Let  $\mathcal{P}^k$  define the space of polynomials of degree  $k$  on element  $E_j$ . A typical finite element subspace (broken space) is defined as

$$\mathcal{D}_k(\mathcal{E}_h) \equiv \{v : v|_{E_j} \in \mathcal{P}^k(E_j), \forall j = 1, \dots, N_h\}$$



- Define the “broken” norm for positive integer  $m$

$$|||\phi|||_{H^m} \equiv \left( \sum_{j=1}^{N_h} \|\phi\|_{H^m(E_j)}^2 \right)^{1/2}$$

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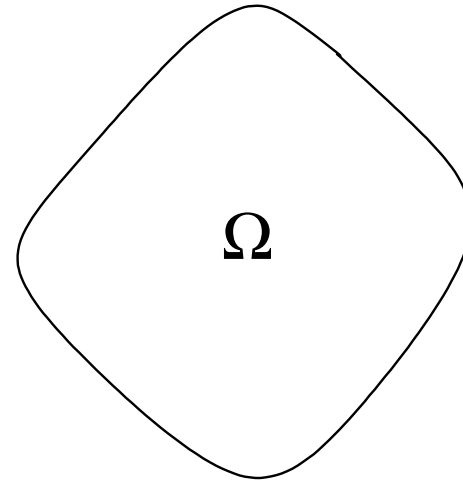
## A Simple Example

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Consider the model problem on bounded domain

$$\begin{aligned} -\Delta u &= f && \text{in } \Omega \\ u &= u_D && \text{on } \partial\Omega_D \\ \nabla u \cdot n &= g \cdot n && \text{on } \partial\Omega_N \end{aligned}$$

We assume that  $f \in L^2(\Omega)$  and  $u_D \in H^{3/2}(\partial\Omega_D)$  and  $g \in H^{1/2}(\partial\Omega_N)$ . Then, there exists a unique solution  $u \in H^1(\Omega)$ .



Weak variational formulations use integration by parts (Green's equation)

$$-\int_R \Delta u = \int_R \nabla u \cdot \nabla v - \int_{\partial R} \nabla u \cdot n v$$

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## Example: continuous Galerkin

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### Weak formulation

Find  $u \in H_0^1(\Omega)$  such that for all  $v \in H_0^1(\Omega)$

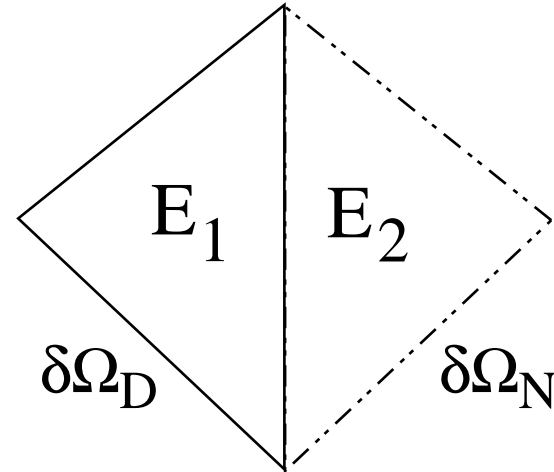
$$\int_{\Omega} \nabla u \nabla v = \int_{\Omega} f v + \int_{\partial\Omega_N} g \cdot n v$$

### Discrete weak formulation

Find  $u_h \in W_h \subset H_0^1(\Omega)$  such that for all  $v \in W_h$

$$\int_{\Omega} \nabla u_h \nabla v = \int_{\Omega} f v + \int_{\partial\Omega_N} g \cdot n v$$

$$W_h(\Omega) \subset H_0^1(\Omega) \equiv \{v \in H^1(\Omega) : v = 0 \text{ on } \partial\Omega\}$$



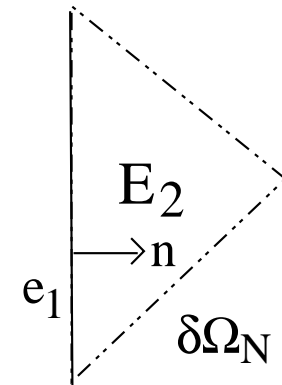
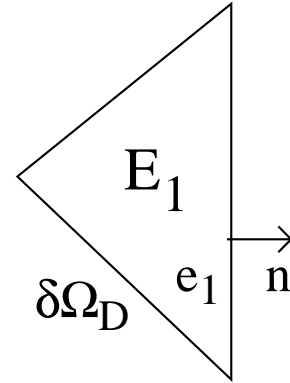
## Example: discontinuous Galerkin

### Weak formulation

Find  $u \in H^1(E_1) \cup H^1(E_2)$  such that for all  $v \in H^1(E_1) \cup H^1(E_2)$

$$\int_{E_1} \nabla u \nabla v - \int_{e_1} \nabla u \cdot n v - \int_{\partial\Omega_D} \nabla u \cdot n v = \int_{E_1} f v$$

$$\int_{E_2} \nabla u \nabla v - \int_{e_1} \nabla u \cdot (-n) v = \int_{E_2} f v + \int_{\partial\Omega_N} g \cdot n v$$

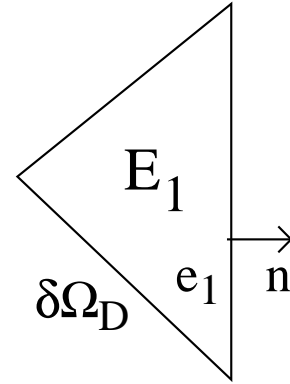


## Example: discontinuous Galerkin

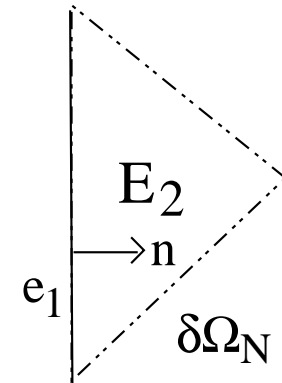
### Discrete weak formulation

Find  $u_h \in \mathcal{D}(\mathcal{E}_h)$  such that for all  $v \in \mathcal{D}(\mathcal{E}_h)$

$$\int_{E_1} \nabla u_h \nabla v - \int_{e_1} \nabla u_h^- \cdot n v^- - \int_{\partial\Omega_D} \nabla u_h \cdot n v = \int_{E_1} f v$$



$$\int_{E_2} \nabla u_h \nabla v + \int_{e_1} \nabla u_h^+ \cdot n v^+ = \int_{E_2} f v + \int_{\partial\Omega_N} g \cdot n v$$



$$\mathcal{D}_k(\mathcal{E}_h) \equiv \{v : v|_{E_j} \in \mathcal{P}^k(E_j), \forall j = 1, \dots, N_h\}$$

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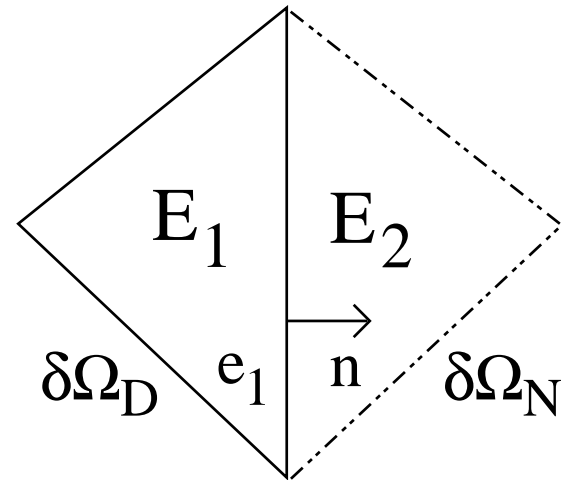
## Example: discontinuous Galerkin

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### Discrete weak formulation

Find  $u_h \in \mathcal{D}(\mathcal{E}_h)$  such that for all  $v \in \mathcal{D}(\mathcal{E}_h)$

$$\begin{aligned} & \sum_{j=1}^2 \int_{E_j} \nabla u_h \nabla v - \int_{e_1} [\nabla u_h \cdot n v] \\ & - \int_{\partial\Omega_D} \nabla u_h^- \cdot n v^- = \sum_{j=1}^2 \int_{E_j} f v + \int_{\partial\Omega_N} g \cdot n v \end{aligned}$$



where the 'jump' is defined as  $[w] = w^- - w^+$  for traces

$$v^-(x) = \lim_{\substack{(s \rightarrow 0^-) \\ (k \rightarrow \infty)}} \gamma_o v_k(x + sn_e), \quad v^+(x) = \lim_{\substack{(s \rightarrow 0^+) \\ (k \rightarrow \infty)}} \gamma_o v_k(x + sn_e),$$



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## Some Issues

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- How to incorporate dirichlet boundary conditions?
- What to do with jump on boundary term (i.e. choose fluxes)?
- Are these methods consistent? Stable? Convergent?
- How efficient are they to implement?
- How accurate are their solutions?
- How do these methods compare with other methods?
- Is there a unified framework for all DG methods?

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## Historical development

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### Elliptic Equations

- First introduced by Nitsche in 1971, application to second-order elliptic equation
- Also known as (interior) penalty methods, attributed to Wheeler in 1978
- Vigorous development in past 5 years

### Hyperbolic Equations

- First introduced by Reed and Hill in 1973, application to neutron transport problem
- First a priori analysis was published by LeSaint and Raviart in 1974
- Vigorous development in past 15 years (including Jaffré)

J. Nitsche, “Über ein Variationsprinzip zur Lösung von Dirichlet Problemen bei Verwendung von Teilräumen, die keinen Randbedingungen unterworfen sind”, *Abh. Math. Sem. Univ. Hamburg*, 36:9-15, 1971

I. Babuška, “The finite element method with interior penalty”, *Math. Comp.*, 27:221-228, 1973

W. Reed and T. Hill, “Triangular mesh methods for the neutron transport equation”, *Los Alamos Scientific Laboratory*, LA-UR-73-479, Los Alamos, NM, 1973

P. LeSaint and P. Raviart, “On a finite element method for solving the neutron transport equations”, in *Mathematical Aspects of Finite Elements in Partial Differential Equations*, Academic Press, 1974

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## DG on elliptic: Nitsche's method

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Consider

$$\begin{aligned} -\Delta u &= f && \text{in } \Omega && \int_{\Omega} \nabla u \cdot \nabla v - \int_{\partial\Omega} \nabla u \cdot n v = \int_{\Omega} f v \\ u &= 0 && \text{on } \partial\Omega \end{aligned}$$

Since  $u$  vanishes on the boundary,

$$\int_{\Omega} \nabla u \cdot \nabla v - \int_{\partial\Omega} \nabla u \cdot n v - \int_{\partial\Omega} \nabla v \cdot n u - \int_{\partial\Omega} \eta u v = \int_{\Omega} f v$$

Find  $u_h \in W_h \subset H^1(\Omega)$  such that for all  $v \in W_h$

$$\int_{\Omega} \nabla u_h \cdot \nabla v - \int_{\partial\Omega} \nabla u_h \cdot n v - \int_{\partial\Omega} \nabla v \cdot n u_h - \int_{\partial\Omega} \eta u_h v = \int_{\Omega} f v$$

consistency term

symmetry term

boundary penalty term

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## Interior Penalties

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Mary Fanett Wheeler

Graduated 1971 Rice University (Houston, TX)

Professor at University of Texas at Austin

Likes scottie dogs and George W. Bush

Idea attributed to private communication by J. Douglas and T. Dupont.  
Summarized and analyzed by D. Arnold in 1979 thesis.

- Suppose approximation space is discontinuous, then

$$\sum_{j=1}^M \int_{E_j} \nabla u \cdot \nabla v - \sum_{i=1}^m \int_{e_i} [\nabla u \cdot n_i v] - \int_{\partial\Omega} \nabla u \cdot n v = \sum_{j=1}^M \int_{E_j} f v$$

M. Wheeler, “An elliptic collocation-finite element method with interior penalties”, *SIAM J. Numer. Anal.*, 15:152-161, 1978

D. Arnold, “An interior penalty finite element method with discontinuous elements”, *SIAM JNA*, 19(4):742-760, 1982

J. Douglas and T. Dupont, “Interior penalty procedures for elliptic and parabolic Galerkin methods”, *Lecture Notes in Physics*, 58, Springer-Verlag, Berlin, 1976.

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## DG Formulation

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- Trick: Define average  $w = \frac{1}{2}(w^- + w^+)$  and recall jump  $[w] = w^- - w^+$  then

$$[a b] = [a]\{b\} + \{a\}[b]$$

- Apply to formulation

$$\begin{aligned} \sum_{j=1}^M \int_{E_j} \nabla u \cdot \nabla v - \sum_{i=1}^m \int_{e_i} [\nabla u \cdot n_i] \{v\} - \sum_{i=1}^m \int_{e_i} \{\nabla u \cdot n_i\} [v] \\ - \int_{\partial\Omega} \nabla u \cdot n v = \sum_{j=1}^M \int_{E_j} f v \end{aligned}$$

- However, since  $u \in H^2(\Omega)$  and  $v \in H^1(\Omega)$ ,

$$\begin{aligned} - \sum_{i=1}^m \int_{e_i} [\nabla u \cdot n_i] \{v\} &= 0, \\ - \sum_{i=1}^m \int_{e_i} \{\nabla v \cdot n_i\} [u] &= 0 \end{aligned}$$

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## Symmetric Interior Penalty DG

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- Idea: add *interior penalty* term of the form

$$\sum_{i=1}^m \int_{e_i} \frac{\sigma}{|e_i|} [u] [v] = 0$$

$$\begin{aligned} \sum_{j=1}^M \int_{E_j} \nabla u \cdot \nabla v - \sum_{i=1}^m \int_{e_i} \{\nabla u \cdot n_i\} [v] - \sum_{i=1}^m \int_{e_i} \{\nabla v \cdot n_i\} [u] \\ + \sum_{i=1}^m \int_{e_i} \frac{\sigma}{|e_i|} [u] [v] - \int_{\partial\Omega} \nabla u \cdot n v = \sum_{j=1}^M \int_{E_j} f v \end{aligned}$$

## Discrete Symmetric Interior Penalty DG

- Can weakly incorporate dirichlet boundary data  $u = g$  on  $\partial\Omega$

$$- \int_{\partial\Omega} (\nabla v \cdot n) u_h = - \int_{\partial\Omega} (\nabla v \cdot n) g$$

Find  $u_h \in \mathcal{D}(\mathcal{E}_h)$  such that for all  $v \in \mathcal{D}(\mathcal{E}_h)$

$$\begin{aligned} \sum_{j=1}^M \int_{E_j} \nabla u_h \cdot \nabla v - \sum_{i=1}^m \int_{e_i} \{\nabla u_h \cdot n_i\} [v] - \sum_{i=1}^m \int_{e_i} \{\nabla v \cdot n_i\} [u_h] \\ + \sum_{i=1}^m \int_{e_i} \frac{\sigma}{|e_i|} [u_h] [v] - \int_{\partial\Omega} (\nabla u_h \cdot n) v - \int_{\partial\Omega} (\nabla v \cdot n) u_h \\ = \sum_{j=1}^M \int_{E_j} f v - \int_{\partial\Omega} \nabla v \cdot n g \end{aligned}$$

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## Non-Symmetric DG Method

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J. Tinsley Oden

Graduated from Oklahoma State University in 1962

Director of TICAM at University of Texas at Austin

Likes hunting big game and has a ranch in Texas



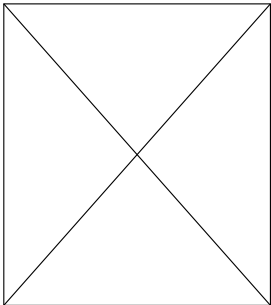
Ivo Babuška

Graduated from Technical University, Prague in 1951

Professor at University of Texas at Austin

Has 4 PhD degrees and an arrest warrant.

Likes studying the history of mathematicians.



Carlos Baumann

Graduated from University of Texas in 1997

Researcher at Computational Mechanics, Inc in Austin

Oden's student at TICAM



## Non-Symmetric Discontinuous Galerkin Method

Find  $u_h \in \mathcal{D}(\mathcal{E}_h)$  such that for all  $v \in \mathcal{D}(\mathcal{E}_h)$

$$\begin{aligned} \sum_{j=1}^M \int_{E_j} \nabla u_h \cdot \nabla v - \sum_{i=1}^m \int_{e_i} \{\nabla u_h \cdot n_i\} [v] + \sum_{i=1}^m \int_{e_i} \{\nabla v \cdot n_i\} [u_h] \\ - \int_{\partial\Omega} (\nabla u_h \cdot n) v + \int_{\partial\Omega} (\nabla v \cdot n) u_h \\ = \sum_{j=1}^M \int_{E_j} f v + \int_{\partial\Omega} \nabla v \cdot n g \end{aligned}$$

- Form becomes positive-definite  $\rightarrow$  inf-sup condition holds  $\rightarrow$  unique solution and conditions of stability
- Consistent, locally mass-conservative formulation

C. Baumann, “An h-p adaptive discontinuous finite element method for computational fluid dynamics”, *PhD thesis*, The University of Texas at Austin, 1997.

J.T. Oden, I. Babuška, and C. Baumann, “A discontinuous hp finite element method for diffusion problems”, *Journ. Comput. Phys.*, 146:491-519, 1998.

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## Non-Symmetric Interior Penalty DG Method

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Béatrice Rivière

Graduated 2000 from University of Texas at Austin  
Assistant Professor at University of Pittsburg  
From Reunion Island, likes to eat Dodo au Vin



Mary Wheeler  
(see page 12)



Vivette Girault  
Professor at Paris VI  
Visited University of Texas at Austin  
Likes math a lot!

## Non-Symmetric Interior Penalty DG Method

Find  $u_h \in \mathcal{D}(\mathcal{E}_h)$  such that for all  $v \in \mathcal{D}(\mathcal{E}_h)$

$$\begin{aligned}
 \sum_{j=1}^M \int_{E_j} \nabla u_h \cdot \nabla v &- \sum_{i=1}^m \int_{e_i} \{\nabla u_h \cdot n_i\} [v] + \sum_{i=1}^m \int_{e_i} \{\nabla v \cdot n_i\} [u_h] \\
 &- \int_{\partial\Omega} (\nabla u_h \cdot n) v + \int_{\partial\Omega} (\nabla v \cdot n) u_h \\
 &+ \sum_{i=1}^m \int_{e_i} \frac{\sigma}{|e_i|} [u_h][v] + \int_{\partial\Omega} \frac{\sigma}{|e_i|} u_h v \\
 &= \sum_{j=1}^M \int_{E_j} f v + \int_{\partial\Omega} \frac{\sigma}{|e_i|} g v + \int_{\partial\Omega} \nabla v \cdot n g
 \end{aligned}$$

- No longer locally mass-conservative, proved convergence rates.

B. Rivière, “Discontinuous Galerkin methods for solving the miscible displacement problem in porous media”, *PhD Thesis*, The University of Texas at Austin, 2002.

B. Rivière, M. Wheeler, and V. Girault, “Improved energy estimates for interior penalty, constrained and discontinuous Galerkin methods for elliptic problems, Part I”, *Comput. Geo*, 8:337-360, 1999.

B. Rivière, M. Wheeler, and V. Girault, “A priori error estimates for finite element methods based on discontinuous approximation spaces for elliptic problems”, *SIAM JNA*, 39(3):902-931, 2001.

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## DG on hyperbolic: Reed and Hill

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- Consider the (linear) neutron transport equation for real number  $\eta$  and constant vector  $a$

$$\eta u + \nabla \cdot (a u) = f \quad \text{in } \Omega$$

Multiply by test function, integrate by parts over element  $E$

$$\int_E \eta u v - \int_E u a \cdot \nabla v - \int_{\partial E} a \cdot n u^\uparrow = \int_E f v \quad \text{in } \Omega$$

where upwind value  $u^\uparrow = \lim_{s \downarrow 0} (x - sa)$  is the value of  $u$  upstream in the characteristic direction  $a$ .

- LeSaint and Raviart immediately publish convergence results (1974), prove (suboptimal)  $h^k$  convergence in  $L^2$  for general triangulations.
- Further results, extensions and numerics are published.

B. Cockburn, E. Karniadakis, C.W. Shu, “The Development of discontinuous Galerkin methods”, In *Discontinuous Galerkin Methods*, Lecture Notes in Computational Science and Engineering, Springer, Berlin, 2000

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## DG on nonlinear hyperbolic systems

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$$\int_E u_t v - \sum_{i=1}^d \int_E f_i(u) \cdot \nabla v + \sum_{i=1}^d \int_{\partial E} f_i(u) \cdot n v = 0$$

- *Time discretization*

Problem: nonlinearities in flux function on boundary prevents element-by-element computation using characteristics

implicit time discretization very inefficient

space-time element construction

- *Analysis*

Problem: nonlinear analysis extremely difficult

only 3 results exist



Jerome Jaffré

Graduated from Paris VI in 1974

Researcher at INRIA Rocquencourt

Likes geosciences a lot!

J. Jaffré, C. Johnson, A. Szeppessy, “Convergence of the discontinuous Galerkin finite element method for hyperbolic conservation laws”, *Math. Models and Meth. in Appl. Sci.*, 5:367-386, 1995

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## Development of LDG method

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- *Cockburn & Shu*: RKDG methods
  - discretize in space via DG, polynomials of degree  $k$
  - discretize in time via explicit TVD Runge-Kutta ( $k + 1$ )
  - employ a general slope limiter to enforce maximum principle
- *Bassi & Rebay*: extension to Navier-Stokes
  - convection-diffusion type equations
  - rewrite as first order system
- *Cockburn & Shu*: LDG methods
  - extension to convection-diffusion systems
  - local solution (parallelizable)

F. Bassi and S. Rebay, “A high-order accurate discontinuous finite element method for the numerical solution of the compressible Navier-Stokes equations”, *J. Comput. Phys.*, 131:267-279, 1997

B. Cockburn and C.W. Shu, “The local discontinuous Galerkin method for time-dependent convection-diffusion systems”, *SIAM JNA*, 35(6):2440-2463, 1998

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## LDG method on transport

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Clint Dawson

Graduated 1988 from Rice University

Professor at University of Texas at Austin

Grew up on a farm in rural Texas

- *Cockburn & Dawson*

- extension to multi-dimension transport equation
- diffusion tensor depends on  $(\mathbf{x}, t)$  and allows for non-invertible diffusion tensors

$$\begin{aligned}\frac{\partial c}{\partial t} + \nabla \cdot (\mathbf{u}c) - \nabla \cdot (D\nabla c) &= 0 & (\mathbf{x}, t) \in \Omega, t > 0, \\ (\mathbf{u}c - D\nabla c) \cdot \mathbf{n} &= (\mathbf{u} \cdot \mathbf{n})\hat{g} & \text{on } \partial\Omega_i \\ (-D\nabla c) \cdot \mathbf{n} &= 0 & \text{on } \partial\Omega_o\end{aligned}$$

B. Cockburn and C. Dawson, “Some extensions of the local discontinuous Galerkin method for convection-diffusion equations in multidimensions”,

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## LDG method on convection-diffusion

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$$\begin{aligned} \frac{\partial c}{\partial t} + \nabla \cdot (\mathbf{u}c - D\nabla c) &= 0 & (\mathbf{x}, t) \in \Omega, t > 0, \\ (\mathbf{u}c - D\nabla c) \cdot \mathbf{n} &= (\mathbf{u} \cdot \mathbf{n})\hat{g} & \text{on } \partial\Omega_i \\ (-D\nabla c) \cdot \mathbf{n} &= 0 & \text{on } \partial\Omega_o \end{aligned}$$

Rewrite as first order system (let  $D=\text{constant}$ )

$$\begin{aligned} \frac{\partial c}{\partial t} + \nabla \cdot (\mathbf{u}c - \nabla \cdot (D\nabla c)) &= 0 \\ \mathbf{z} &= -D\nabla c \\ (\mathbf{u}c - \mathbf{z}) \cdot \mathbf{n} &= (\mathbf{u} \cdot \mathbf{n})\hat{g} & \text{on } \partial\Omega_i \\ \mathbf{z} \cdot \mathbf{n} &= 0 & \text{on } \partial\Omega_o \end{aligned}$$



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## LDG method

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Apply DG to each equation separately

$$\int_E \frac{\partial c}{\partial t} v - \int_E (\mathbf{u}c + \mathbf{z}) \cdot \nabla v + \int_{\partial E} (\mathbf{u}c + \mathbf{z}) \cdot \mathbf{n} v = \int_E f v$$

$$\int_E \mathbf{z} \mathbf{w} - \int_E Dc \nabla \cdot \mathbf{w} + \int_{\partial E} Dc \mathbf{w} \cdot \mathbf{n} = 0$$

$$\int_{\partial E \cap \partial \Omega_i} (\mathbf{u}c - z) \cdot \mathbf{n} v = \int_{\partial E \cap \partial \Omega_i} (\mathbf{u} \cdot \mathbf{n}) \hat{g} v$$
$$\int_{\partial E \cap \partial \Omega_o} \mathbf{u}c \cdot \mathbf{n} v = 0$$

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## Discrete LDG method

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$$\begin{aligned}
 - \sum_{j=1}^M \int_{E_j} \frac{\partial c_h}{\partial t} v - \sum_{j=1}^M \int_{E_j} (\mathbf{u}c_h + \mathbf{z}_h) \cdot \nabla v + \sum_{i=1}^m \int_{e_i} (\mathbf{u}c_h^\uparrow + \{\mathbf{z}_h\}) \cdot \mathbf{n} [v] \\
 + \int_{\partial\Omega_o \cap \partial\Omega} (\mathbf{u}c_h \cdot \mathbf{n}) v + \int_{\partial\Omega_i \cap \partial\Omega} \hat{g}\mathbf{u} \cdot \mathbf{n} v = \sum_{j=1}^M \int_{E_j} f v
 \end{aligned}$$

$$\sum_{j=1}^M \int_{E_j} \mathbf{z}_h \cdot \mathbf{w}_h + \sum_{j=1}^M \int_{E_j} Dc_h \nabla \cdot \mathbf{w}_h - \sum_{i=1}^m \int_{e_i} \{Dc_h\} [\mathbf{w}_h] \cdot \mathbf{n}_i = 0$$

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## Features

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### DISCONTINUOUS GALERKIN FEM

- easily incorporates h- and p-adaptivity
- locally conservative and accurate
- easily handles complex geometry and nonmatching grids
- suppresses non-physical oscillations
- handles advection dominated flow regimes efficiently
- degrees of freedom associated with elements

### CONTINUOUS GALERKIN FEM

- handles diffusion dominated flow regimes efficiently
- degrees of freedom usually associated with nodes