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# Long time stability of Markov and branching processes: a pedagogical introduction.

INRIA – Séminaire des doctorants

Grégoire FERRÉ, Gabriel STOLTZ, Mathias ROUSSET

CERMICS – ENPC, INRIA & LABEX BEZOUT

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1. What are we talking about ? What for ?
2. Markov chains and their stability
3. Branching processes

1. What are we talking about ? What for ?

# A classical description of matter

In classical mechanics, a common way of describing a system at a fixed temperature  $T$  is to consider:

- a **set of atoms** with positions  $x$ ;
- a **potential energy** given by a Hamiltonien  $H(x)$ ;
- a **force**  $F$  acting on the system, typically typically  $F(x) = -\nabla H(x)$ ;
- a thermal agitation at temperature  $T$ , generally represented by a **random noise**.

# Mathematical formulation

In practice, the latter elements are resumed in the following time evolution, for  $n \in \mathbb{N}$ :

$$x_{n+1} = x_n + F(x_n) + \sigma G_n,$$

where:

- $n$  corresponds to time;
- $F(x_n)$  is the force at location  $x_n$ ;
- $\sigma$  is the intensity of the noise, proportional to  $\sqrt{T}$ ;
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When  $F = -\nabla H$ , we have

$$x_{n+1} = x_n - \nabla H(x_n) + \sigma G_n,$$

hence  $x_{n+1}$  «reduces» the energy  $H$  in average.

# Some motivations and questions

This basic framework can be used to model many different systems:

- molecules, chemical reactions;
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- surface interactions, etc.

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The dynamics of the system is encoded in **the transition kernel  $P$** , for a position  $x$  and a set  $S$ ,

$$P(x, S) = \mathbb{P}\left[x_{n+1} \in S \mid x_n = x\right],$$

that is **the probability of reaching a set  $S$  starting from  $x$** . There are **natural questions** about such systems:

- their **long time behavior** (Sec. 2. );
- the **probabilities of fluctuations** (Sec. 3. );
- decorrelation times, linear response theory, etc.



## 2. Markov chains and their stability

# What is stability?

Importance of the **long time behavior** for the computation of averages of functions  $\varphi$  (like a pressure), as:

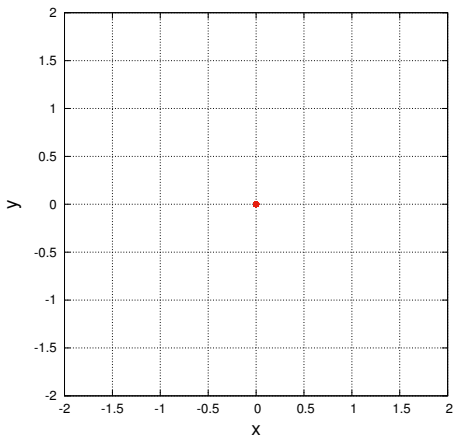
$$\frac{1}{N_{\text{iter}}} \sum_{n=0}^{N_{\text{iter}}-1} \varphi(x_n) \xrightarrow{N_{\text{iter}} \rightarrow +\infty} \text{average of } \varphi.$$

For this, an important condition to hold is **stability**. By this, we mean somehow that the **system does not get lost at infinity**. Two useful applications:

- more **physical insight** into the considered system;
- stability of **numerical schemes**.

# A basic example

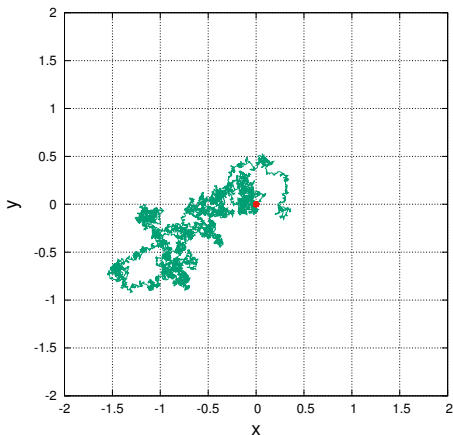
Brownian motion over  $\mathbb{R}^2$ :  $F = 0$ , so  $(x_n)$  is a **sum of  $2d$ -Gaussian** random variables.



The process starts at  $(0,0)$ ...

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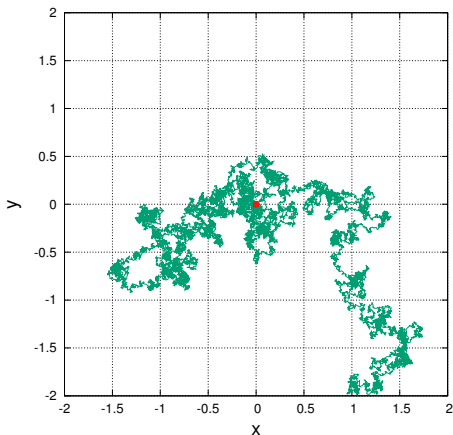
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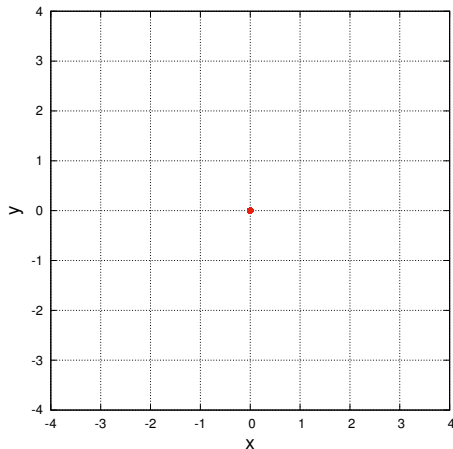
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and goes away without  
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# Partially confined particle

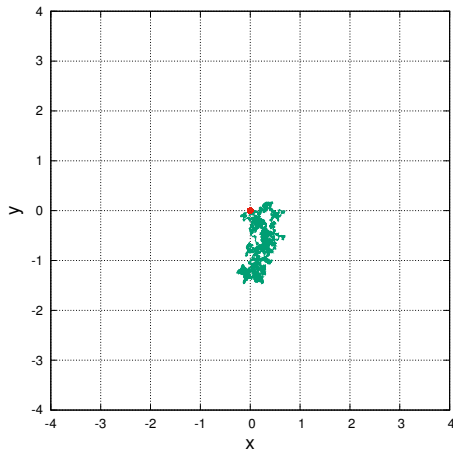
Dynamics over  $\mathbb{R}^2$  with  $F(x; y) = (-5x; 0)$ .



The process is confined in the  $x$  direction...

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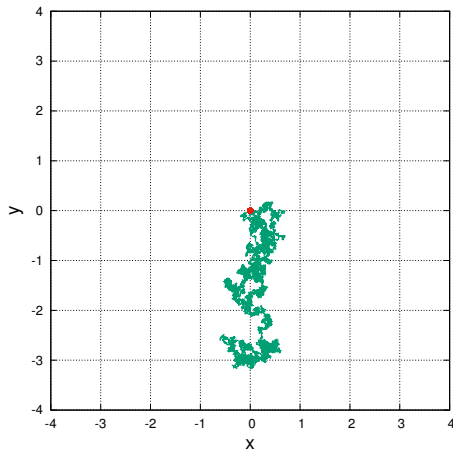
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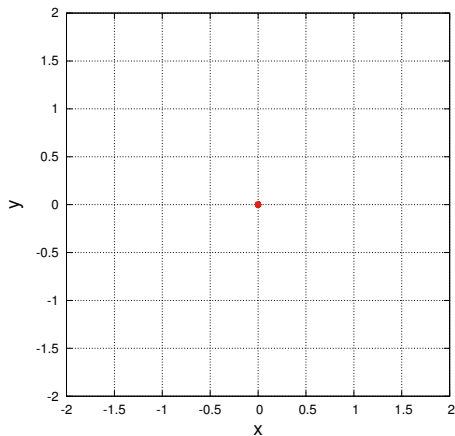


The process is confined in the  $x$  direction... but moves without restriction in the  $y$  direction.



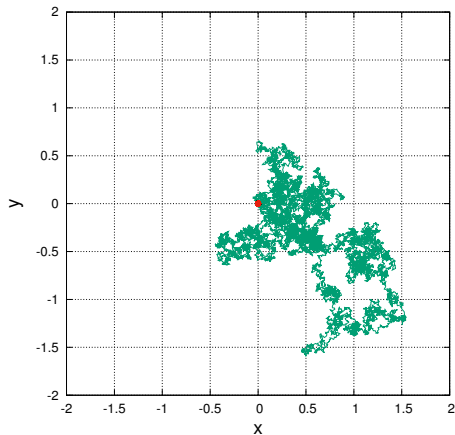
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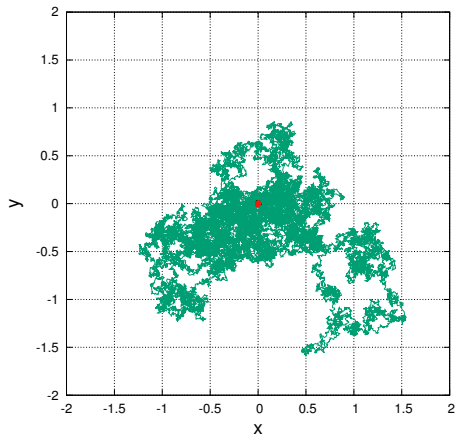
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# Hand waving theorem

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A possible solution: show the **existence of an energy** that **decreases in average**.

## Idea of theorem

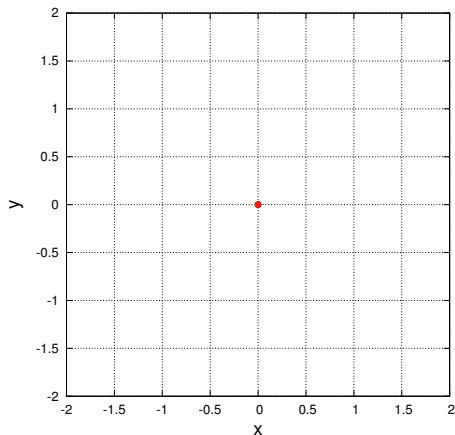
If there exists a function  $W \geq 0$  and constants  $\gamma \in (0, 1)$ ,  $b \geq 0$  such that

$$PW(x) \leq \gamma W(x) + b,$$

and  $W(x) \rightarrow +\infty$  as  $|x| \rightarrow +\infty$ , then the dynamics is stable.

# Back to the Ornstein-Uhlenbeck process

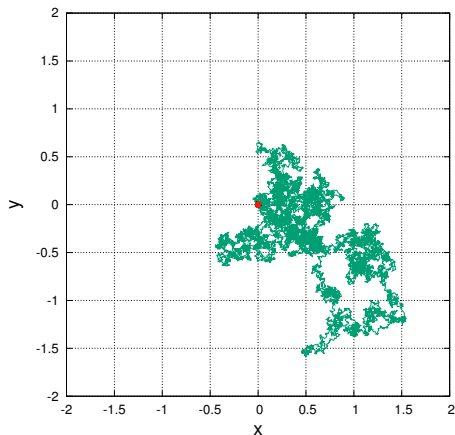
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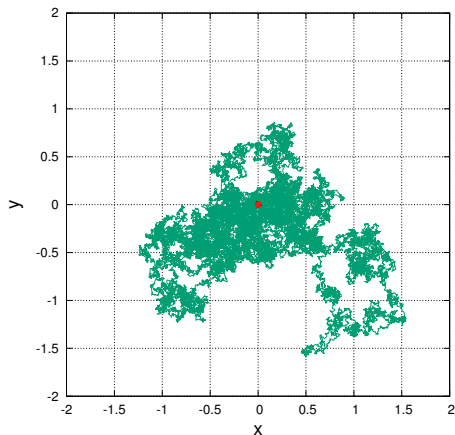
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The force  $F = -\nabla H$  derives from the energy

$$H(x, y) = x^2 + y^2.$$

In practice, one can choose

$$W(x, y) = H(x, y),$$

the energy of the system.



### 3. Branching processes

# What is a branching process?

A branching system consists of:

- a set of  $M$  replicas  $(x^m)_{m=1}^M$ ;
- each replica follows a Markovian dynamics;
- each replica  $m$  is assigned a weight  $w_m$  depending on its past trajectory;
- the total weight is  $\bar{w} = \sum_{m=1}^M w_m$ ;
- one computes a probability of surviving  $p_m = w_m/\bar{w}$ ;
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In practice, the weight over  $k$  steps of the  $m^{\text{th}}$  replica is computed with

$$w_m^k = \exp \left( \sum_{i=0}^{k-1} V(x_i^m) \right).$$

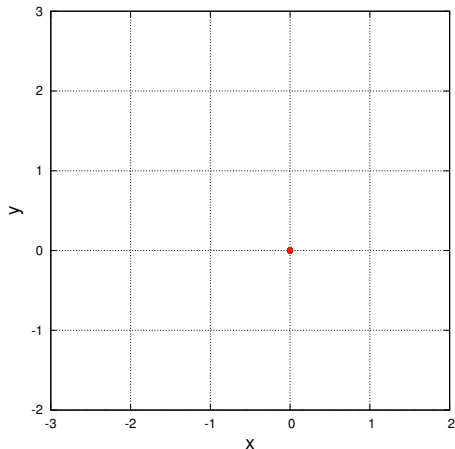
# Some motivations

There are (at least) three motivations for studying such dynamics:

- in quantum physics;
- in rare event simulation;
- in filtering.

# A quantum mechanics example

Brownian dynamics  $(x_n)$  over  $\mathbb{R}^2$  with  $V(x) = -|x|^2$ .



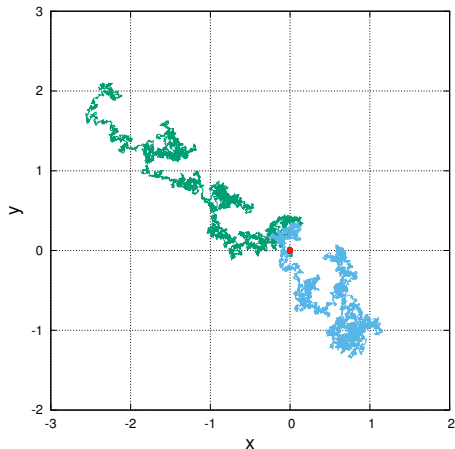
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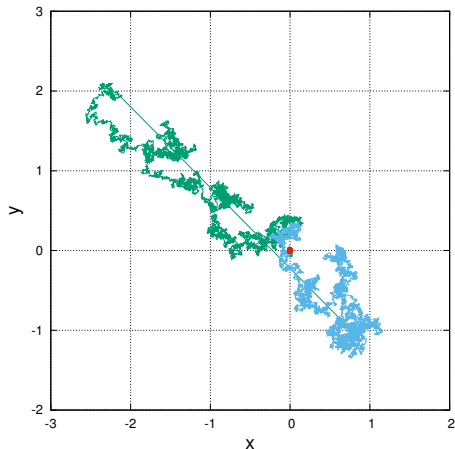
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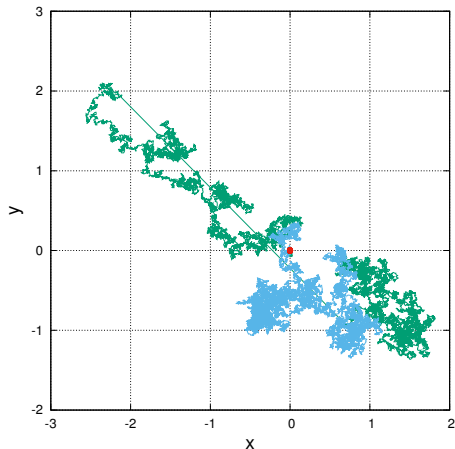
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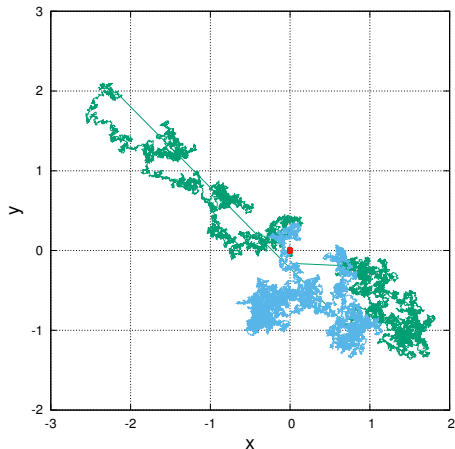
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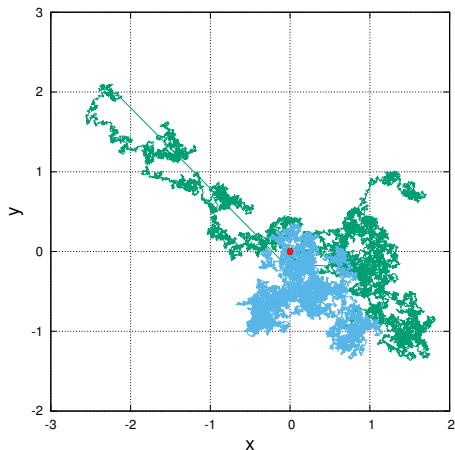
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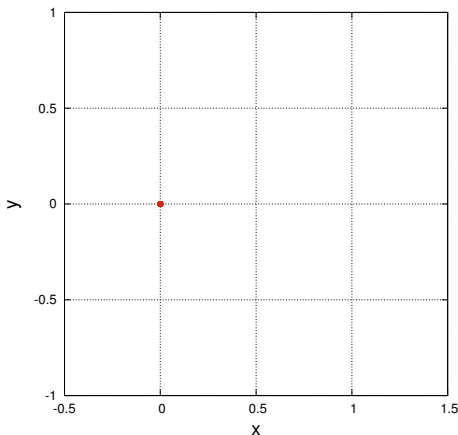
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**Conclusion:** there is a confinement by selection.

# A rare event example

Ornstein-Uhlenbeck dynamics with  $F(x,y) = (-2x, -2y)$ . Weight function  $V(x,y) = x$ .



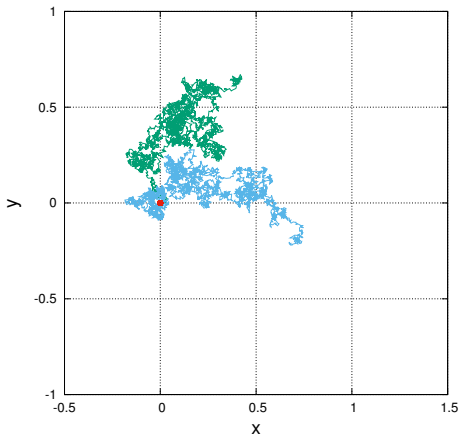
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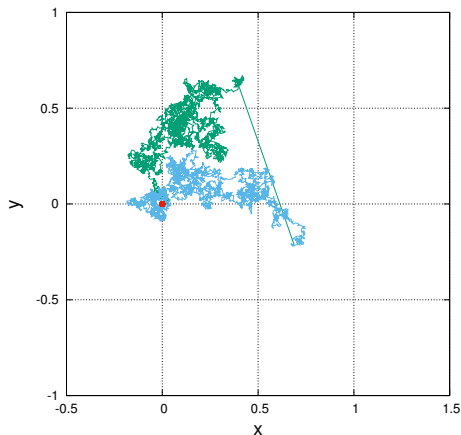
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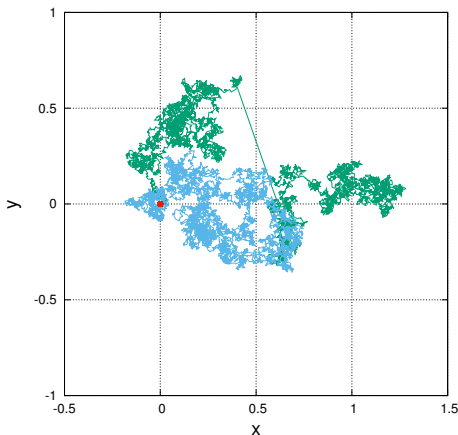
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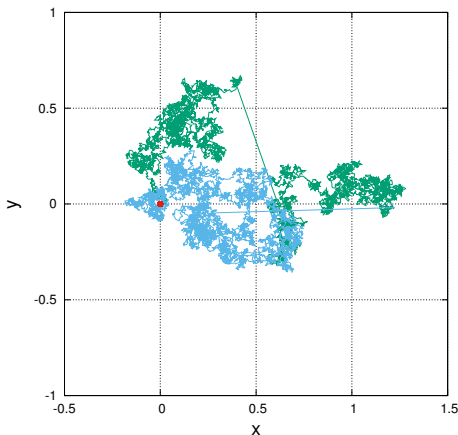
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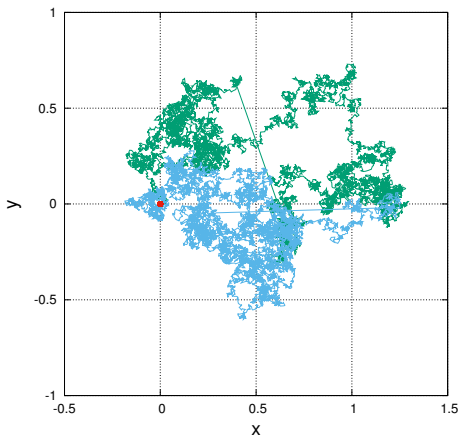
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**Conclusion:** the particles are selected towards the right.



# Some heuristic

First observations:

- there are two phenomena: the **dynamics** itself and the **selection** rule;
- the confinement can be provided by the dynamics or by the selection;

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Two natural questions:

- **unstable** Markov dynamics (first example)  $\rightarrow$  **stable** branching process ?
- **stable** Markov dynamics (stable example)  $\rightarrow$  **unstable** branching process ?

# Hand waving new theorem

Condition of stability based on the **decrease of some energy** ? The answer is yes [G.F., M. Rousset, G. Stoltz, in prep.]. For this, define

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**Question: what the ❄ does that mean?**

# Hand waving of hand waving etc

**Idea:** the energy decrease must take into account the weight  $V$ .

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## Result in a nutshell

More generally,

- if  $V(x) \rightarrow -\infty$  as  $|x| \rightarrow +\infty$ , the dynamics is **always stable**, since replicas always end up dying;
- if not, there is a trade-off: too many kids away may **blow up to your face**.

Take home message:

- random systems for modeling physics;
- problems of stability, long time behavior and energy;
- branching processes, a **funny problem** (as far as mathematics can be funny);
- new results, «energy» for branching processes.